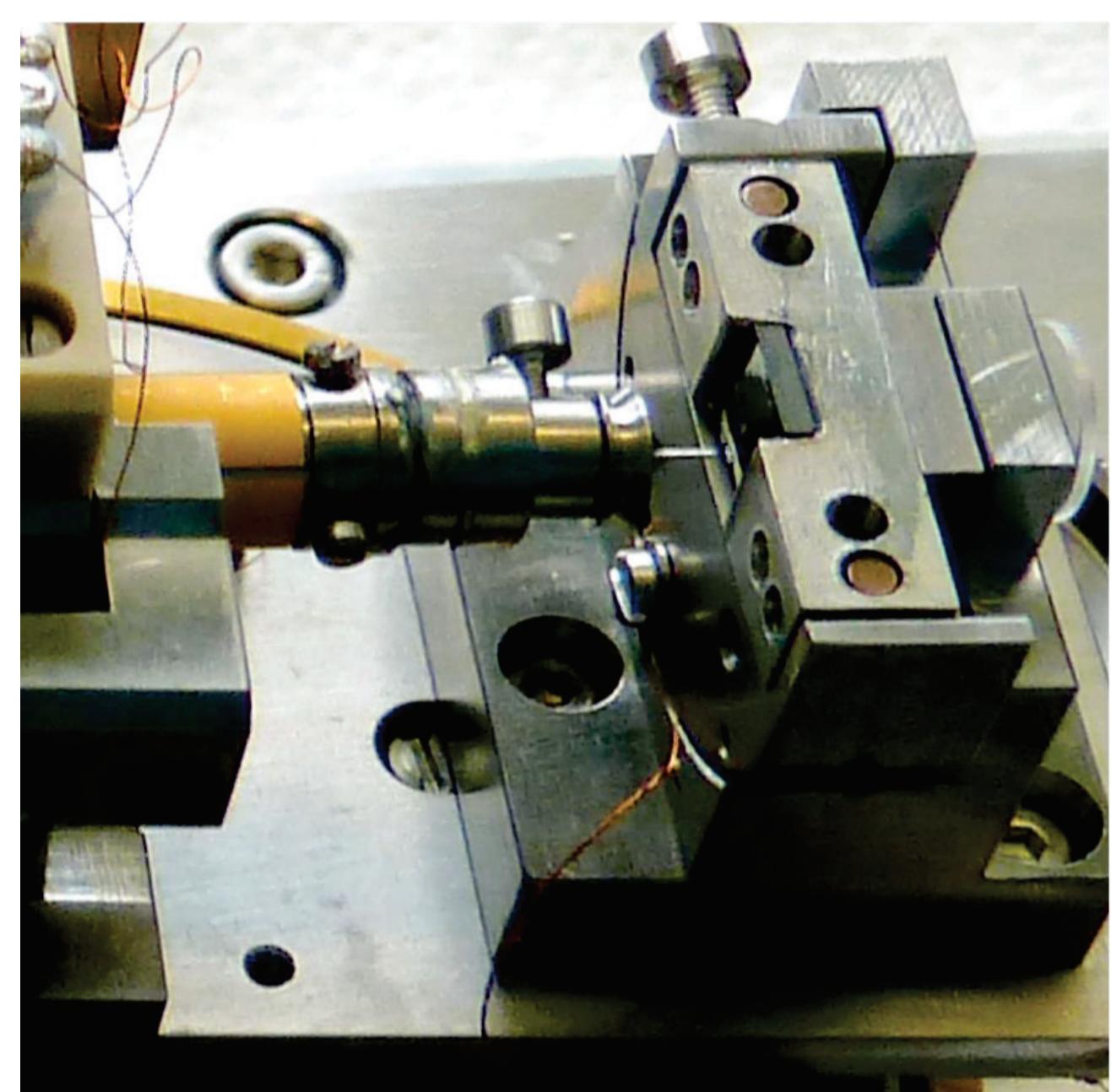


# Looking for the Origin of Power Laws in Electric Field Assisted Tunneling

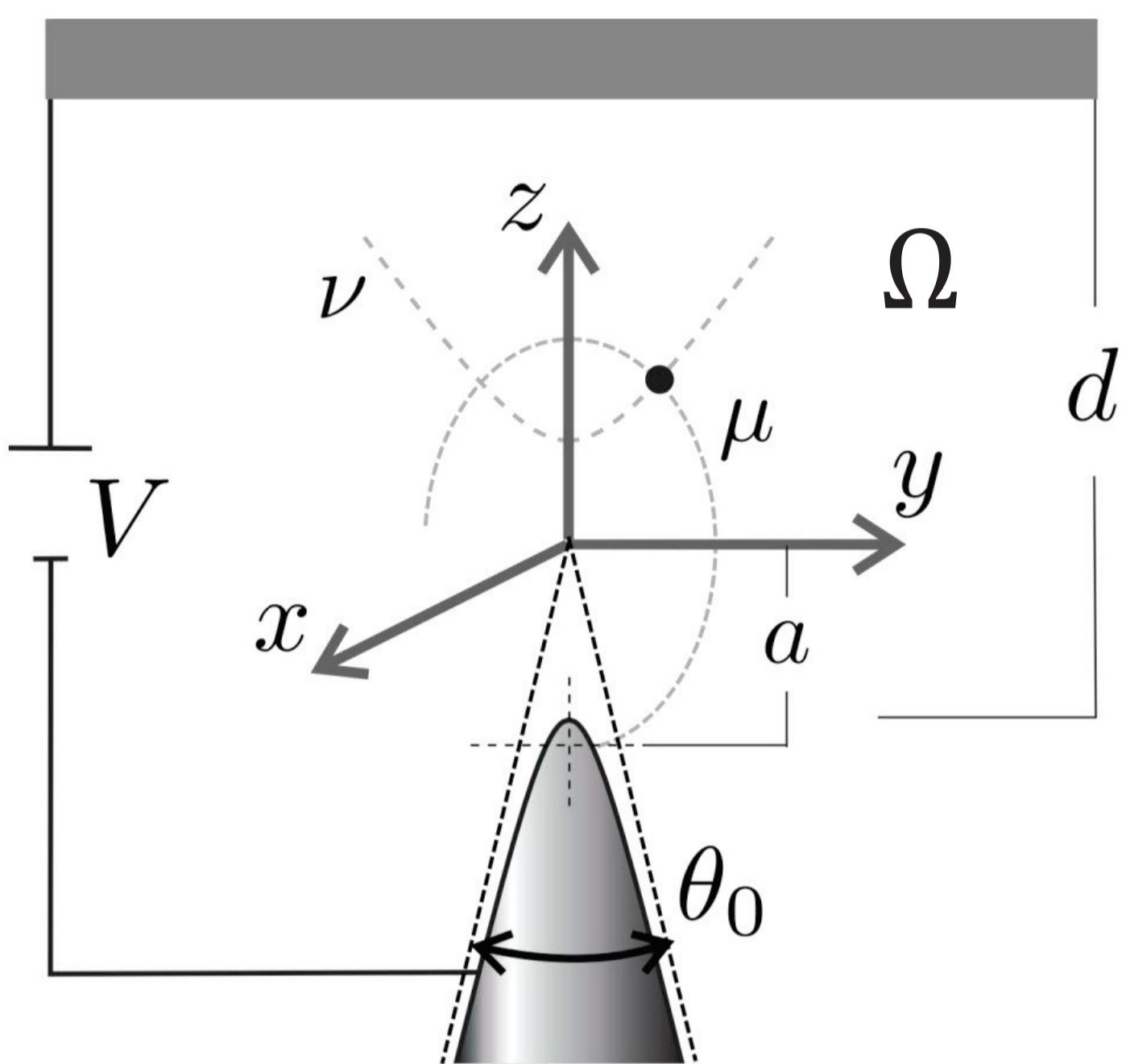
H. Cabrera, D.A. Zanin, L.G. De Pietro, A. Vindigni, U. Ramsperger, D. Pescia

Laboratory for Solid State Physics, ETH Zurich, Wolfgang-Pauli-Strasse 16, Zurich 8093, Switzerland.

**Introduction:** A sharp tip approached perpendicular to a conducting surface at subnanometer distances and biased with a small voltage builds a junction across which electrons can be transferred from the tip apex to the nearest surface atom by direct quantum mechanical tunneling. Such a junction is used e.g. in Scanning Tunneling Microscopy (STM).



**Figure 1.** The experimental Setup: a nanoelectronic device



**Figure 2.** The electrostatic problem, tip modeled with a hyperboloid of revolution

When the distance  $d$  between tip and collector is increased beyond some nanometers, the junction enters the electric field assisted regime, the one underlying the topografiner technology –an imaging technique widely used in micro- and nano-electronics. Recent experiments<sup>1</sup> in this regime suggest a scaling law which can be tested numerically by verifying the collapsing of a family of  $\Phi(z, d)$ -curves, computed at different  $d$ , onto one single curve ( $\Phi(z, d)$  being the electric potential).

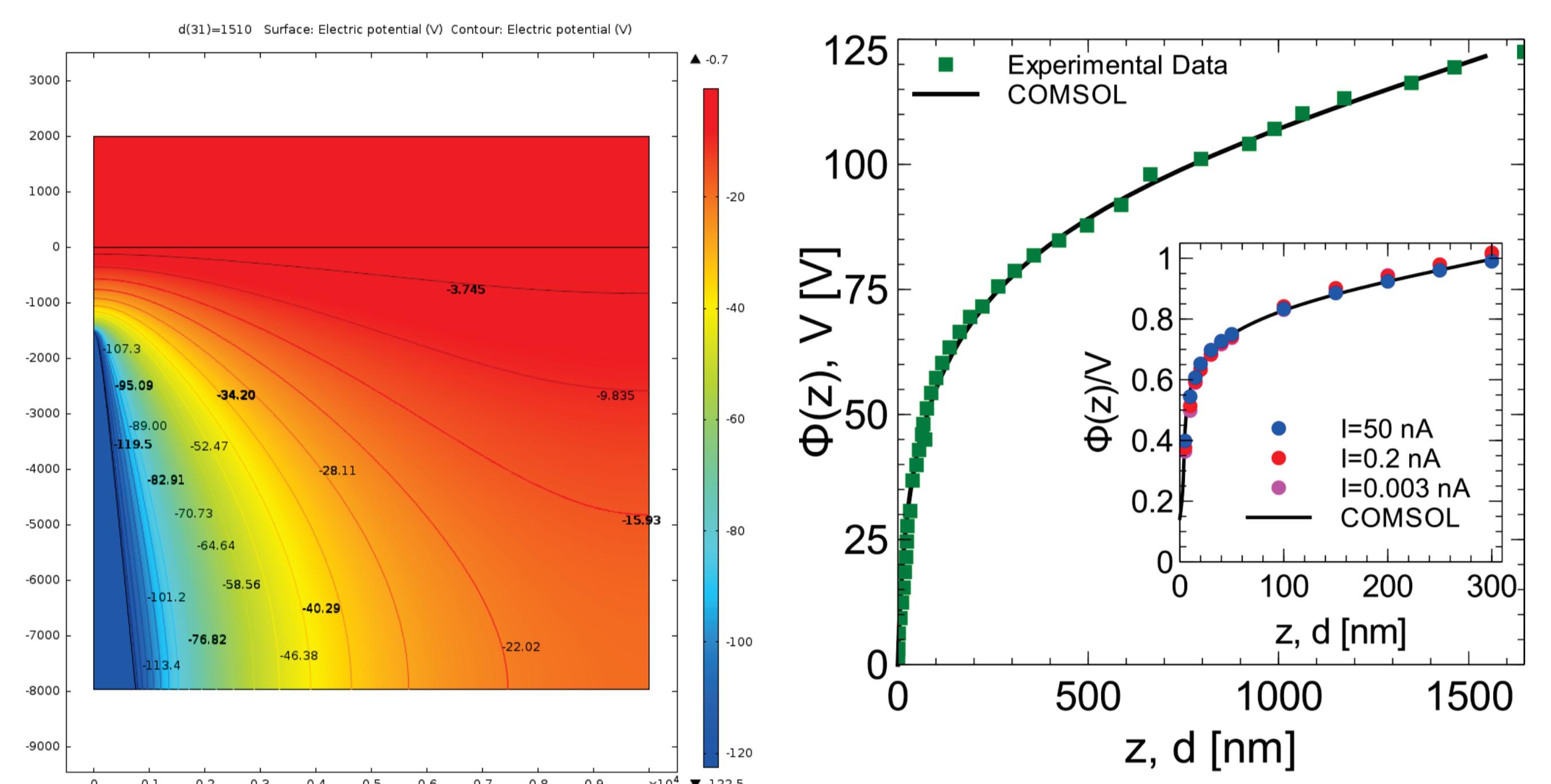
**Computational Methods:** The tip, kept at ground potential, is placed in front of a conducting plane set at voltage  $+V$  (see Fig. 2). Denoting with  $\Omega$  the region of space excluding the tip and the plane, the electrostatic problem is a well defined Dirichlet problem for the electric potential  $\Phi$ :

$$\begin{cases} \nabla^2 \Phi = 0 & \text{in } \Omega, \\ \Phi = 0 & \text{on the surface of the tip,} \\ \Phi = +V & \text{on the plane,} \\ |\Phi|(x) \leq V & \forall \vec{r} = (x, y, z) \in \Omega. \end{cases}$$

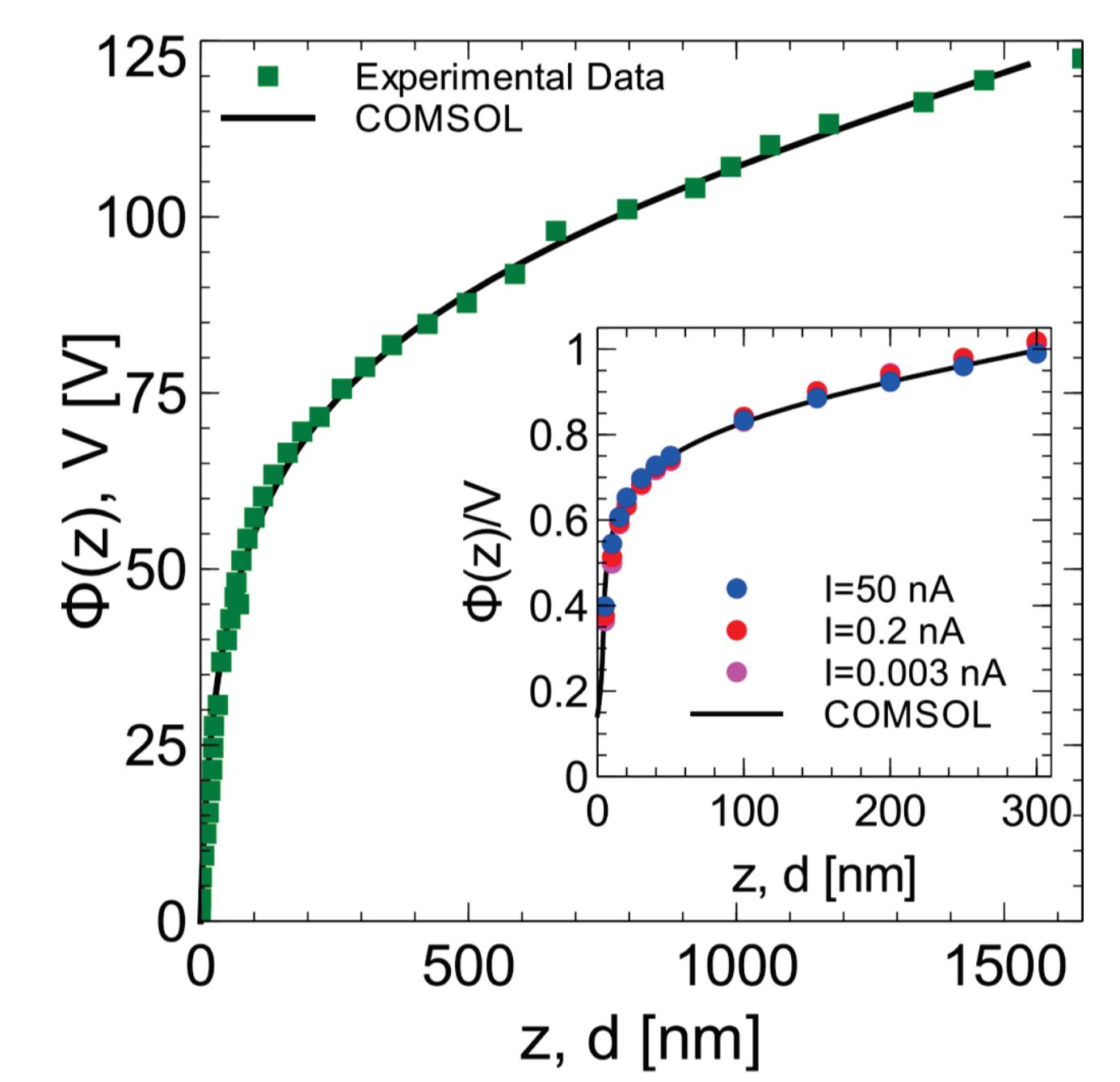
**Results:** To the leading order, the potential depends on  $d$  differently in the “near” ( $d \ll a$ ) or “distant” ( $d \gg a$ ) regime

$$\begin{cases} \Phi(z) \sim V \cdot \frac{z}{d} & \text{for } d \ll a, \\ \Phi(z) \sim V \cdot \left(\frac{a}{d}\right)^\lambda \cdot \frac{z}{a} & \text{for } d \gg a, \end{cases}$$

with the geometry-factor  $\lambda$  determined by the aperture angle of the tip  $\theta_0$  (see Fig. 2).



**Figure 3.** COMSOL simulation of the diodelike junction



**Figure 4.** Examples of  $\Phi/V$ - $d$  curves for constant current

**Conclusions:** In the range  $d \gg 10$  nm, the experimental data follow a power law  $\propto d^\lambda$ , with  $\lambda = 0.21 \pm 0.02$ . For smaller values of  $d$  the dependence becomes almost linear, indicating that the junction behaves as a plane capacitor at short distances: Direct tunneling typically occurs in this geometry. The essential features observed experimentally are captured by introducing the potential  $\Phi$  computed for “realistic” tips into standard equations for electric field assisted tunneling. This highlights the potential of COMSOL-simulation in the context of field-emitted electron microscopy.

## References:

1. H. Cabrera *et al.*, Phys. Rev. B 87, 115436 (2013)
2. D.A. Zanin *et al.*, Advances in Imaging and Electron Physics bf 170, 227 (2012)